

ECS 455 Chapter 1

Introduction & Review

Radio

1.3 Wireless Channel (Part 1)

{ noise
interference } random, time-varying

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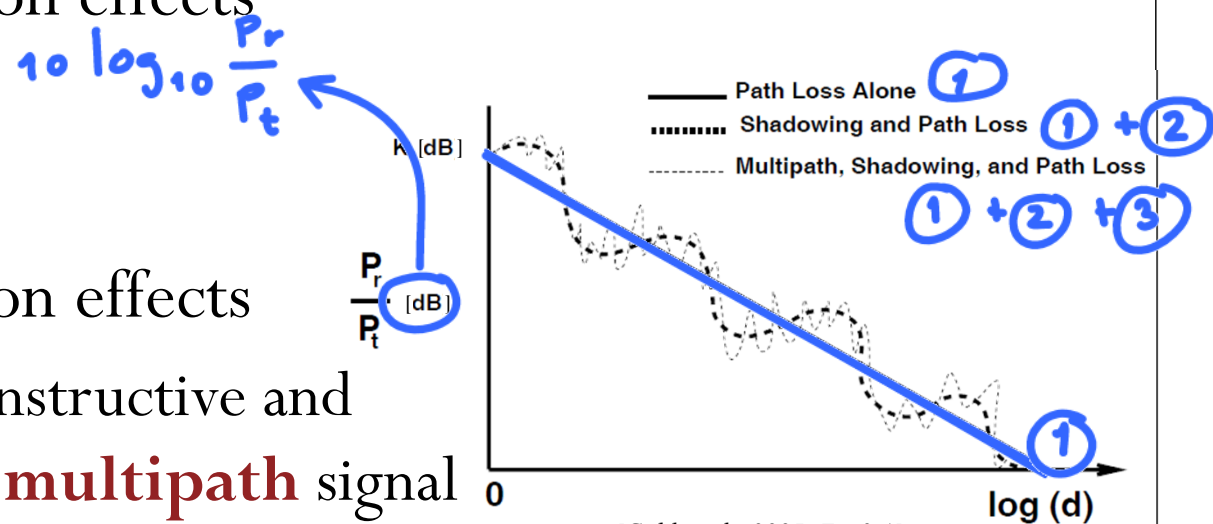
Wireless Channel

- **Large-scale** propagation effects

1. Path loss
2. Shadowing

3. • **Small-scale** propagation effects

- Variation due to the constructive and destructive addition of **multipath** signal components.
- Occur over very **short distances**, on the order of the signal **wavelength**.



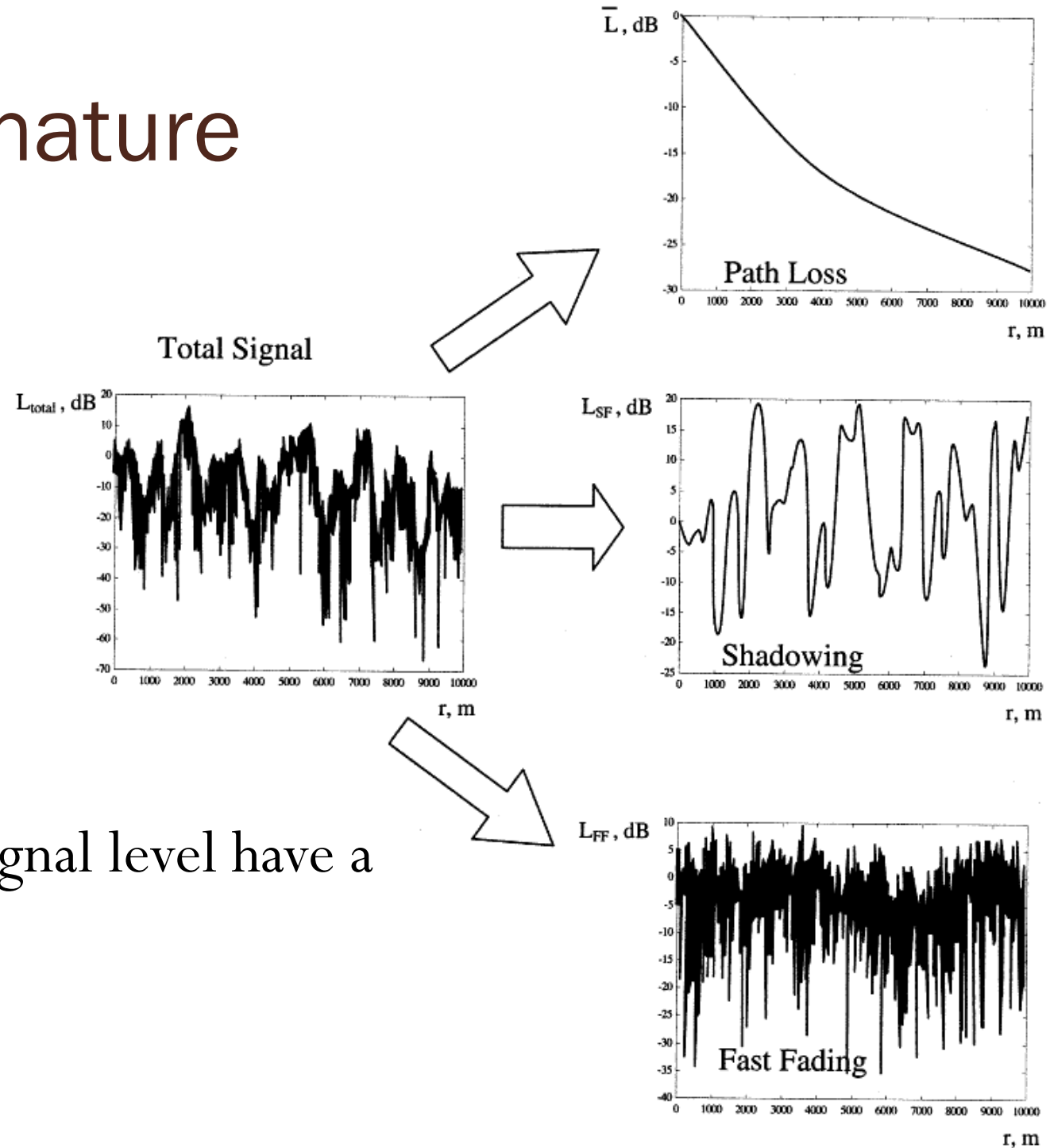
[Goldsmith, 2005, Fig 2.1]

$$\lambda = \frac{c}{f}$$

$c \leftarrow \approx 3 \times 10^8 \text{ [m/s]}$
 $f \rightarrow 3 \times 10^9$

$f = 3 \text{ GHz} \rightarrow \lambda = 0.1 \text{ m}$

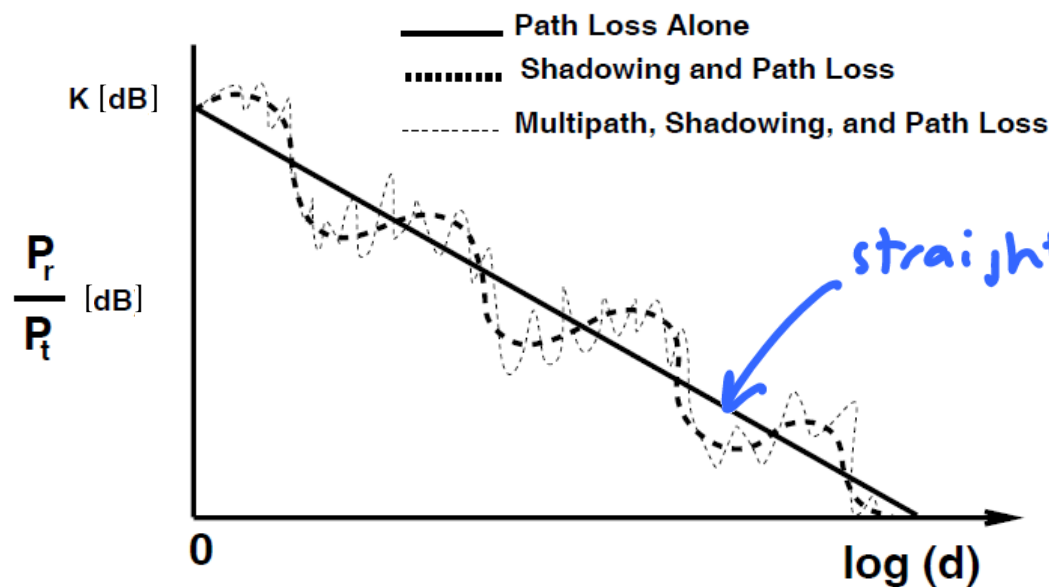
Triple nature



Variations of signal level have a triple nature.

① Path loss

- Caused by
 - **dissipation** of the **power** radiated by the transmitter
 - effects of the propagation channel
- Models generally assume that it is the same at a given transmit-receive distance.
- **Variation occurs over large distances** (100-1000 m)



straight line in log-log plot

[Goldsmith, 2005, Fig 2.1]

Path Loss (PL)

$$P_L = \frac{\text{Transmitted power}}{\text{Average received power}} = \frac{P_t}{P_r}$$

Averaged over any random variations

- **Free-Space** Path Loss Model:

$$\frac{P_r}{P_t} \propto \frac{1}{d^2}$$

- P_r falls off inversely proportional to the square of the distance d between the Tx and Rx antennas.

- **Simplified** Path Loss Model:

$$\frac{P_r}{P_t} = K \left(\frac{d_0}{d} \right)^\gamma$$

To be discussed

(Path loss of the free-space model)

Friis Equation (Free-Space PL)

- One of the most fundamental equations in antenna theory

1 for non-directional antennas

$$\frac{P_r}{P_t} = \left(\frac{\sqrt{G_{Tx} G_{Rx}} \lambda}{4\pi d} \right)^2 = \left(\frac{\sqrt{G_{Tx} G_{Rx}} c}{4\pi d f} \right)^2 \propto \frac{1}{d^2 f^2}$$

- Lose more power at higher frequencies.

0.7 GHz → 2.4 GHz → 5 GHz → 60 GHz

10.7 dB loss

$$20 \log_{10} \frac{2.4}{0.7}$$

6.4 dB loss

$$20 \log_{10} \frac{5}{2.4}$$

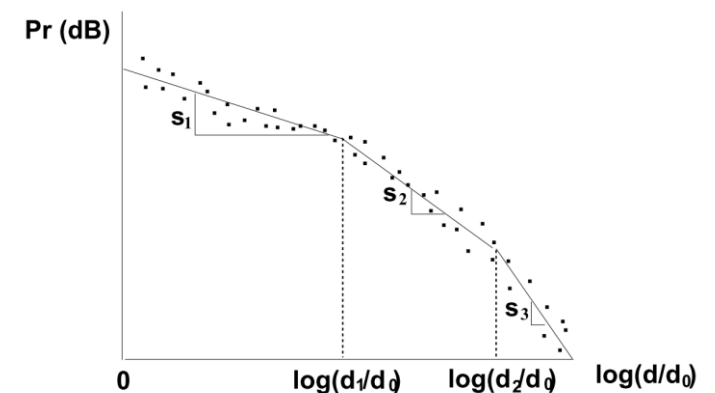
21.6 dB loss

$$20 \log_{10} \frac{60}{5}$$

- Some of these losses can be offset by reducing the maximum operating range.
 - The remaining loss must be compensated for by increasing the antenna gain.

More Path Loss Models

- Analytical models
 - Maxwell's equations
 - Ray tracing
 - Empirical models: Developed to predict path loss in typical environment.
 - Okumura
 - Hata
 - COST 231
 - by EURO-COST (EUROpean COoperative for Scientific and Technical research)
 - Piecewise Linear (Multi-Slope) Model
 - Tradeoff: Simplified Path Loss Model
- Prohibitive (complex, impractical)
Need to know/specify “almost everything” about the environment.



Simplified Path Loss Model

$$\frac{P_r}{P_t} = K \left(\frac{d_0}{d} \right)^\gamma$$

$$10 \log_{10} \frac{P_r}{P_t} \text{ [dB]} = \left(10 \log_{10} K d_0^\gamma \right) - 10\gamma \log_{10} d$$

Captures the essence of signal propagation without resorting to complicated path loss models, which are only approximations to the real channel anyway!

- K is a unitless constant which depends on the antenna characteristics and the average channel attenuation
 - $\left(\frac{\lambda}{4\pi d_0} \right)^2$ for free-space path gain at distance d_0 assuming omnidirectional antennas
- d_0 is a reference distance for the antenna far-field
 - Typically 1-10 m indoors and 10-100 m outdoors.
- γ is the **path loss exponent**.

(Near-field has scattering phenomena.)

Path Loss Exponent γ

- 2 in free-space model
- 4 in two-ray model
[Goldsmith, 2005, eq. 2.17]
- Cellular: 3.5 – 4.5
[Myung and Goodman, 2008 , p 17]
- Larger @ higher freq.
- Lower @ higher antenna heights

Environment	γ range
Urban macrocells	3.7-6.5
Urban microcells	2.7-3.5
Office Building (same floor)	1.6-3.5
Office Building (multiple floors)	2-6
Store	1.8-2.2
Factory	1.6-3.3
Home	3

Indoor Attenuation Factors

- Building penetration loss: 8-20 dB (better if behind windows)
- Attenuation between floors
 - @ 900 MHz
 - 10-20 dB when the Tx and Rx are separated by a single floor
 - 6-10 dB per floor for the next three subsequent floors
 - A few dB per floor for more than four floors
 - Typically worse at higher frequency.
- Attenuation across floors

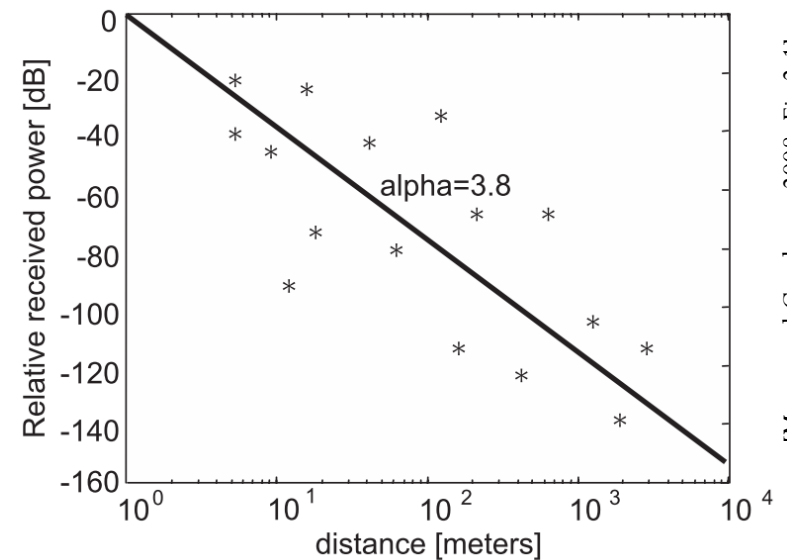
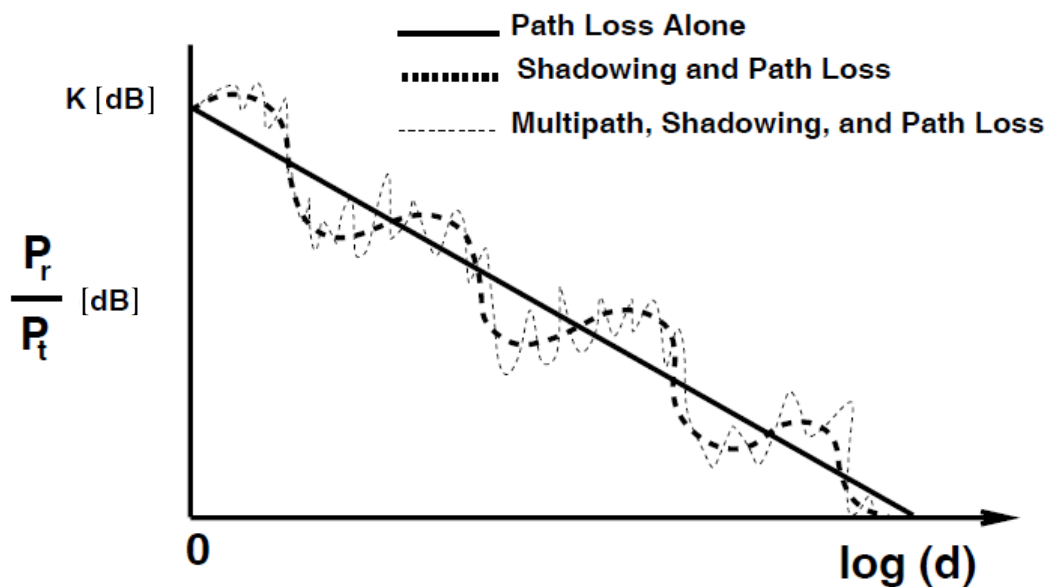
Partition Type	Partition Loss in dB
Cloth Partition	1.4
Double Plasterboard Wall	3.4
Foil Insulation	3.9
Concrete wall	13
Aluminum Siding	20.4
All Metal	26

[Goldsmith, 2005, Sec. 2.5.5]

② Shadowing (or Shadow Fading)

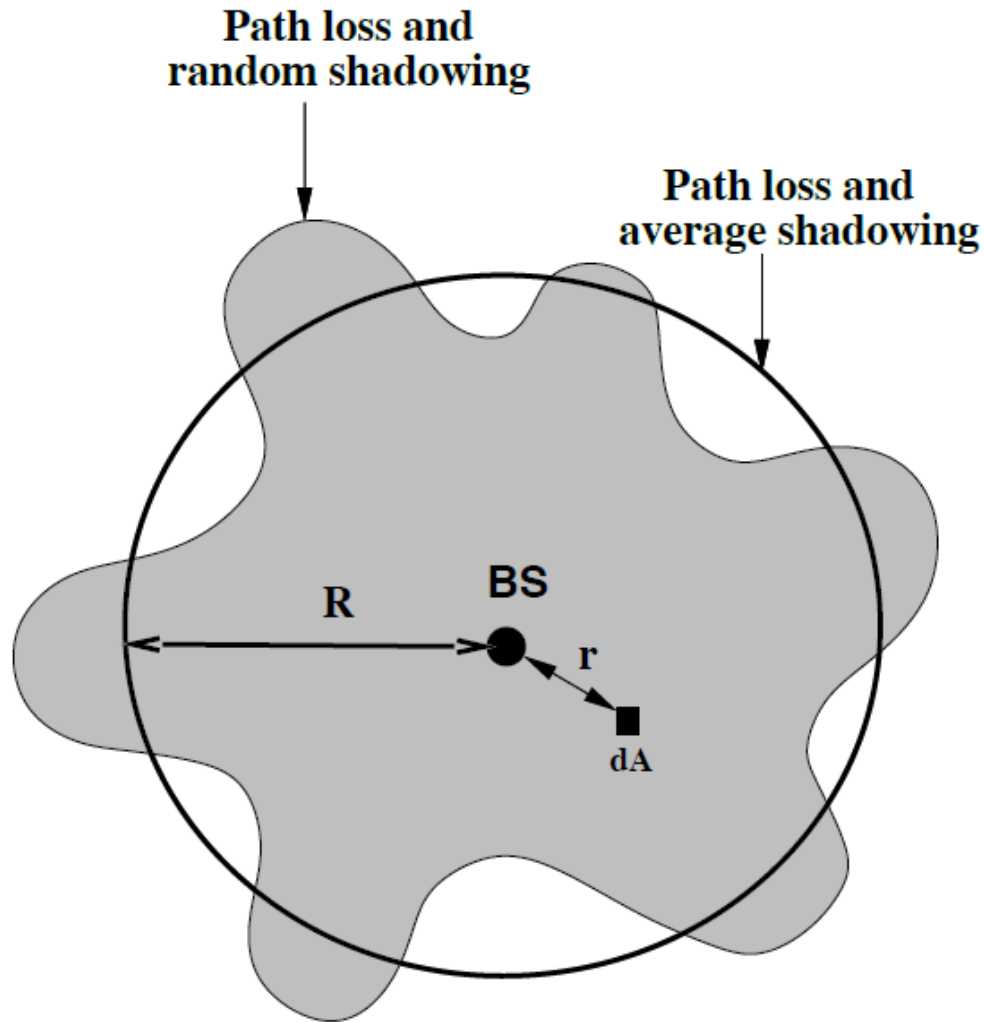
- Additional attenuation caused by **obstacles** (large objects such as buildings and hills) between the transmitter and receiver.
 - Think: cloud blocking sunlight
- Attenuate signal power through absorption, reflection, scattering, and diffraction.
- Variation occurs over distances proportional to the length of the obstructing object (**10-100 m** in outdoor environments and less in indoor environments).

[Goldsmith, 2005, Fig 2.1]



[Myung and Goodman, 2008, Fig 2.1]

Contours of Constant Received Power



[Goldsmith, 2005, Fig 2.10]

Log-normal shadowing

- Random variation due to blockage from objects in the signal path and changes in reflecting surfaces and scattering objects
→ random variations of the received power at a given distance

$$10 \log_{10} \frac{P_t}{P_r} \sim \mathcal{N}(\mu, \sigma^2)$$

4 – 13 dB with higher values in urban areas and lower ones in flat rural environments.

in dB

- This model has been confirmed empirically to accurately model the variation in received power in both outdoor and indoor radio propagation environments.

[Erceg et al, 1999] and [Ghassemzadeh et al, 2003]

Log-normal shadowing (motivation)

- Location, size, dielectric properties of the blocking objects as well as the changes in reflecting surfaces and scattering objects that cause the random attenuation are generally unknown
 \Rightarrow statistical models must be used to characterize this attenuation.
- Assume a large number of shadowing objects between the transmitter and receiver

Without the objects, the attenuation factor is $K \left(\frac{d_0}{d} \right)^\gamma$.

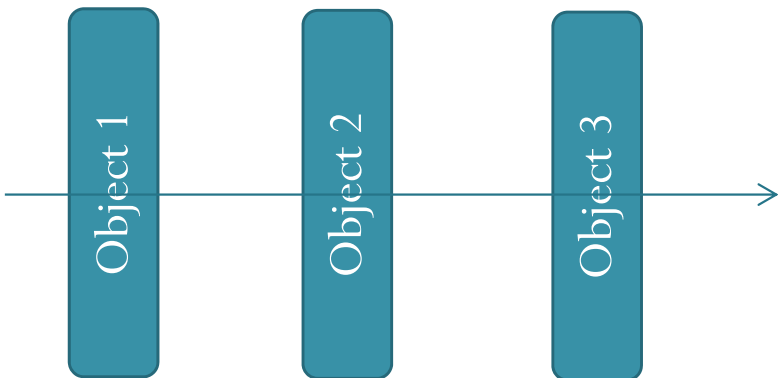
Each object introduces extra power loss factor of α_i .

So,

$$\frac{P_r}{P_t} = K \left(\frac{d_0}{d} \right)^\gamma \prod_i \alpha_i$$

$$10 \log_{10} \frac{P_r}{P_t} = 10 \log_{10} K \left(\frac{d_0}{d} \right)^\gamma - \sum_i 10 \log_{10} \alpha_i$$

$3 + \mathcal{N}(5, 2^2) = \mathcal{N}(\epsilon, \sigma^2)$ By CLT, this is approximately Gaussian



PDF of Lognormal RV

- Consider a random variable

$$R = \frac{P_t}{P_r}$$

- Suppose

$$10 \log_{10} R \sim \mathcal{N}(\mu, \sigma^2)$$

Here, it should be clear why the unit of σ is in dB.

- Then,

$$f_R(r) = \begin{cases} \frac{1}{\sqrt{2\pi}\sigma} \frac{10}{\ln 10} \frac{1}{r} e^{-\frac{1}{2} \left(\frac{(10 \log r) - \mu}{\sigma} \right)^2}, & r > 0 \\ 0, & \text{otherwise.} \end{cases}$$

For typical cellular environment, σ is in the range of 5-12 dB.
[Proakis and Salehi, 2007, p 843]

PDF of Lognormal RV (Proof)

Suppose $c \log_b Y \sim \mathcal{N}(\mu, \sigma^2)$.

Let $X = c \log_b Y$. Note that $X = c \log_b Y = \frac{c}{\ln b} \ln(Y) = k \ln(Y)$.

Then, $Y = e^{\frac{X}{k}}$ where $k = \frac{c}{\ln b}$.

Recall, from ECS315 that to find the pdf of $Y = g(X)$ from the pdf of X , we first find the cdf of Y and then differentiate to get its pdf:

$$F_Y(y) = P[Y \leq y] = P\left[e^{\frac{X}{k}} \leq y\right] = P[X \leq k \ln(y)] = F_X(k \ln(y)).$$

$$f_Y(y) = \frac{d}{dy} F_X(k \ln(y)) = \frac{k}{y} f_X(k \ln(y)) = \frac{1}{\sqrt{2\pi\sigma}} \frac{k}{y} e^{-\frac{1}{2}\left(\frac{k \ln(y) - \mu}{\sigma}\right)^2}.$$

PDF of Lognormal RV (Proof)

Suppose $c \log_b Y \sim \mathcal{N}(\mu, \sigma^2)$.

Let $X = c \log_b Y$. Note that $X = c \log_b Y = \frac{c}{\ln b} \ln(Y) = k \ln(Y)$.

Then, $Y = e^{\frac{X}{k}}$ where $k = \frac{c}{\ln b}$.

Alternatively, to find the pdf of $Y = g(X)$ from the pdf of X , when g is monotone, we may use the formula:

$$f_X(x) |dx| = f_Y(y) |dy| \longrightarrow f_Y(y) = \left| \frac{dx}{dy} \right| f_X(x)$$

This gives $f_Y(y) = \frac{k}{y} f_X(c \log_b y)$ (same as what we found earlier).

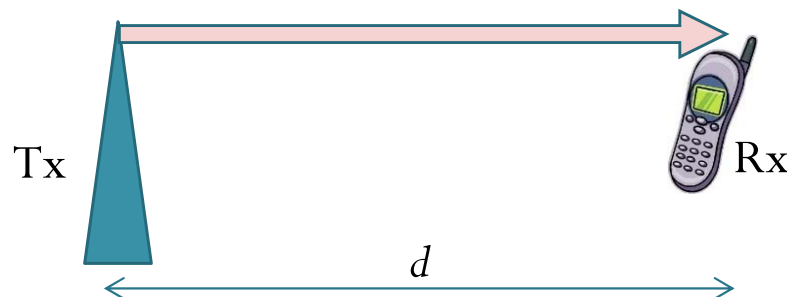
Ray tracing (a prelude)

- Approximate the solution of Maxwell's equations
 - Approximate the propagation of electromagnetic waves by representing the wavefronts as simple **particles**.
 - Thus, the reflection, diffraction, and scattering effects on the wavefront are approximated using **simple geometric equations** instead of Maxwell's more complex wave equations.
- Assumption: the received waveform can be approximated by the sum of the free space wave from the transmitter plus the reflected free space waves from each of the reflecting obstacles.

$$x(t) = \sqrt{2P_t} \cos(2\pi f_c t)$$

$$y(t) = ?$$

$$\cos\left(2\pi f_c \left(t - \frac{d}{c}\right)\right)$$



Energy and Power

$$|z|^2 = z z^*$$

instantaneous power $p(t) = \frac{v^2(t)}{R} = i^2(t) R$

- Consider a signal $g(t)$.
- Total (normalized) **energy**:

$R=1$

Parseval's Theorem

$$E_g = \int_{-\infty}^{\infty} |g(t)|^2 dt = \lim_{T \rightarrow \infty} \int_{-T}^T |g(t)|^2 dt = \int_{-\infty}^{\infty} |G(f)|^2 df$$

$$\langle g, g \rangle = \int_{-\infty}^{\infty} g(t) g^*(t) dt = \int_{-\infty}^{\infty} G(f) G^*(f) df$$

$$\Psi_g(f) = |G(f)|^2$$

ESD: Energy Spectral Density

- Average (normalized) **power**:

$$P_g = \left\langle |g(t)|^2 \right\rangle = \lim_{T \rightarrow \infty} \frac{1}{T} \int_{-T/2}^{T/2} |g(t)|^2 dt = \lim_{T \rightarrow \infty} \frac{1}{2T} \int_{-T}^T |g(t)|^2 dt$$

time average

$$\langle x(t) \rangle = \lim_{T \rightarrow \infty} \frac{1}{2T} \int_{-T}^T x(t) dt$$

Time average operator:

Fact: Suppose $x(t)$ is periodic with period T_0 .

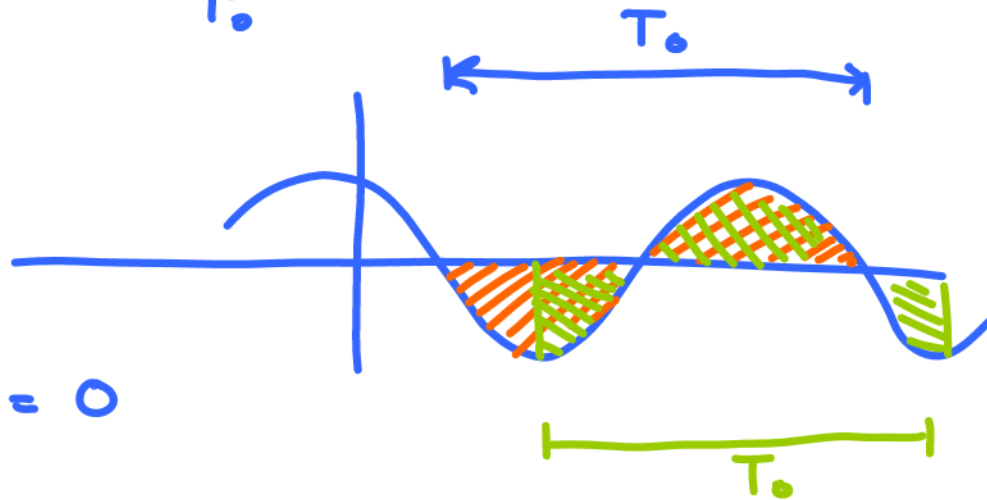
$$\text{Then, } \langle x(t) \rangle = \frac{1}{T_0} \int_{T_0} x(t) dt.$$

Example:

$$\langle \cos(2\pi f_c t + \theta) \rangle = \frac{1}{T_0} \int_{T_0} \cos(2\pi f_c t + \theta) dt = \frac{1}{T_0} \times 0 = 0$$

↓

$$\text{Period } T_0 = \frac{1}{f_c}$$



$$\langle \sin(2\pi f_c t + \theta) \rangle = 0$$

$$\langle \cos^2(2\pi f_c t + \theta) \rangle = \left\langle \frac{1}{2} (1 + \cos(2\pi 2f_c t + 2\theta)) \right\rangle = \frac{1}{2}$$

Three Important Facts

(i) • If $g(t)$ is periodic with period T_0 , then $P_g = \frac{1}{T_0} \int_{T_0} |g(t)|^2 dt$

(ii) • If $g(t)$ can be expressed as $\sum_{k=1}^n c_k e^{j2\pi f_k t}$ where the f_k are distinct, then $P_g = \sum_{k=1}^n |c_k|^2$.

• Ex

$$g(t) = A e^{j(2\pi f_c t + \theta)} \Rightarrow P_g = |A|^2$$

$$g(t) = A \cos(2\pi f_c t + \theta) \Rightarrow P_g = \frac{|A|^2}{2}$$

When $g(t)$ is periodic with period T_0 :

$$g(t) = \sum_{k=-\infty}^{\infty} c_k e^{j2\pi \frac{k}{T_0} t} \Rightarrow P_g = \sum_{k=-\infty}^{\infty} |c_k|^2$$

(iii) • Suppose that $g(t)$ can be expressed as $a(t)\cos(2\pi f_c t + \phi)$ and that $A(f - f_c)$ and $A(f + f_c)$ do not overlap. Then

$$P_g = \frac{1}{2} P_a$$

Ex $g(t) = 1 e^{j2\pi f_0 t} \Rightarrow P_g = 1^2 = 1$

$$P_g = \langle |g(t)|^2 \rangle = \langle |e^{j2\pi f_0 t}|^2 \rangle = \langle 1 \rangle = 1$$

$$g(t) = e^{j2\pi 3t} + e^{j2\pi 4t} \Rightarrow P_g = 1^2 + 1^2 = 2$$

$$g(t) = e^{j2\pi 3t} + e^{j2\pi 3t} \Rightarrow P_g = 2^2 = 4$$

$$= 2 e^{j2\pi 3t} \quad \neq 1^2 + 1^2$$

$$g(t) = A e^{j(2\pi f_c t + \theta)} \Rightarrow P_g = |A e^{j\theta}|^2$$

$$= A e^{j\theta} e^{j2\pi f_c t} \quad = |A|^2$$

$$g(t) = \cos(2\pi f_c t) = \frac{1}{2} (e^{j2\pi f_c t} + e^{-j2\pi f_c t}) \Rightarrow P_g = \left(\frac{1}{2}\right)^2 + \left(\frac{1}{2}\right)^2 = 2 \times \frac{1}{4} = \frac{1}{2}$$

Alternatively, $P_g = \langle \cos^2(2\pi f_c t) \rangle = \frac{1}{2}$

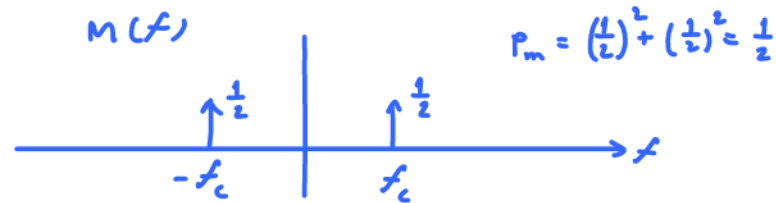
$$g(t) = A \cos(2\pi f_c t + \theta)$$

$$= \frac{A}{2} (e^{j(\pi f_c t + \theta)} + e^{-j(2\pi f_c t + \theta)}) \Rightarrow P_g = \left|\frac{A}{2}\right|^2 + \left|\frac{A}{2}\right|^2 = \frac{|A|^2}{2}$$

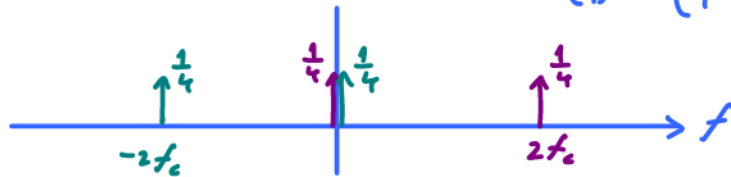
(iii) $x(t) = \underbrace{m(t)}_{\text{shifting}} \sqrt{2} \cos(2\pi f_c t) \Rightarrow P_x = \frac{1}{2} P_{\sqrt{2}m(t)} = \frac{1}{2} \times 2 P_m = P_m$
 shifted copies do not overlap
 $\langle (c m(t))^2 \rangle = \langle c^2 m^2(t) \rangle = c^2 \langle m^2(t) \rangle$

When they overlap:

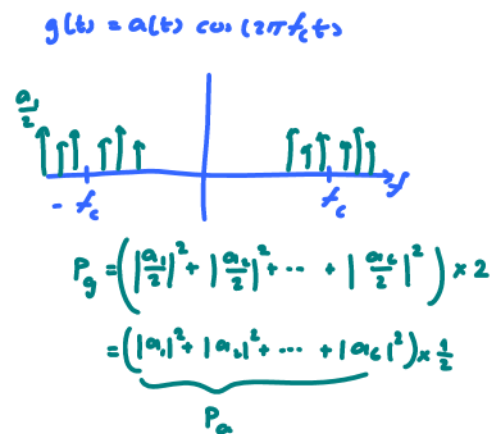
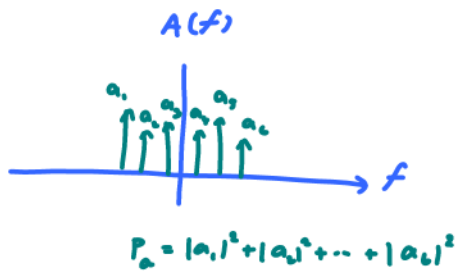
$m(t) = \cos(2\pi f_c t) \Rightarrow P_m = \frac{1}{2}$



$m(t) \cos(2\pi f_c t) \Rightarrow P = \left(\frac{1}{4}\right)^2 + \left(\frac{1}{4} + \frac{1}{4}\right)^2 + \left(\frac{1}{4}\right)^2 = \frac{3}{8} \neq \frac{P_m}{2}$



"Proof" of (iii)



Power Calculation (More Examples)

non-overlapping
 $P_{\Sigma} = \Sigma P$

iv) $g(t) = \sum_k a_k(t) \cos(2\pi f_k t + \phi_k)$

$$P_g = \frac{1}{2} \sum_k P_{a_k} = \sum_k \frac{P_{a_k}}{2}$$

Assume the $A_k(f \pm f_k)$'s do not overlap.

v) $g(t) = a_1 \cos(2\pi f_c t + \phi_1) + a_2 \cos(2\pi f_c t + \phi_2)$

same

$$P_g = \frac{1}{2} |a_1 e^{j\phi_1} + a_2 e^{j\phi_2}|^2$$

$$= \frac{1}{2} a_1^2 + \frac{1}{2} a_2^2 + a_1 a_2 \cos(\phi_2 - \phi_1)$$

Ex. $g(t) = 4 \cos(2t) + 3 \sin(2t)$
 $= 4 \cos(2t) + 3 \cos(2t - 90^\circ)$
 $\equiv 4 \angle 0^\circ + 3 \angle -90^\circ$
 $= 5 \angle -36.87^\circ$
 $= 5 \cos(2t - 36.87^\circ)$

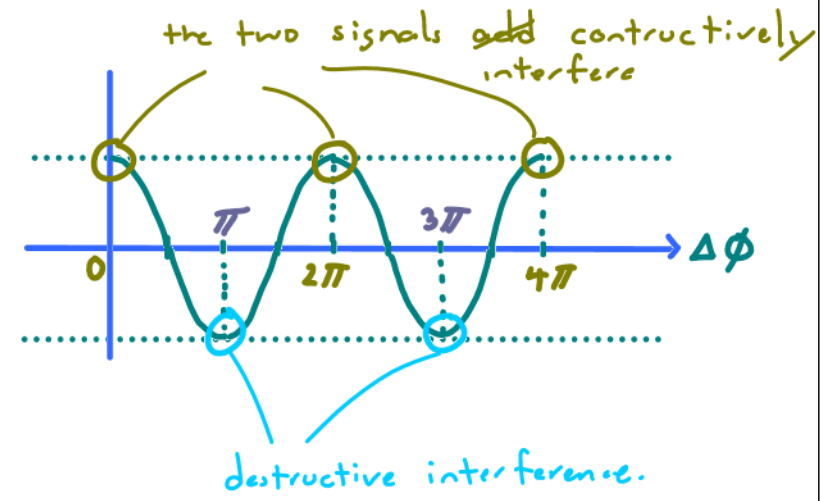


$$P_g = \frac{\Sigma^2}{2}$$

$$(a_1 - a_2)^2 \geq 0$$

$$a_1^2 + a_2^2 - 2a_1 a_2 \geq 0$$

$$\frac{a_1^2}{2} + \frac{a_2^2}{2} \geq a_1 a_2$$



$$\begin{aligned}
 g(t) &= a_1 \cos(2\pi f_c t + \phi_1) + a_2 \cos(2\pi f_c t + \phi_2) \\
 &= \operatorname{Re} \left\{ a_1 e^{j(2\pi f_c t + \phi_1)} \right\} + \operatorname{Re} \left\{ a_2 e^{j(2\pi f_c t + \phi_2)} \right\} \\
 &= \operatorname{Re} \left\{ (a_1 e^{j\phi_1} + a_2 e^{j\phi_2}) e^{j(2\pi f_c t)} \right\} \\
 &= \operatorname{Re} \left\{ a e^{j\phi} e^{j2\pi f_c t} \right\} = \operatorname{Re} \left\{ a e^{j(2\pi f_c t + \phi)} \right\} \\
 &= a \cos(2\pi f_c t + \phi)
 \end{aligned}$$

$$z = a e^{j\phi}$$

$$P_g = \frac{a^2}{2}$$

$$\begin{aligned}
 |z|^2 &= z z^* = (a_1 e^{j\phi_1} + a_2 e^{j\phi_2}) (a_1 e^{-j\phi_1} + a_2 e^{-j\phi_2}) \\
 &= a_1^2 + a_2^2 + a_1 a_2 e^{j(\phi_1 - \phi_2)} + a_1 a_2 e^{-j(\phi_1 - \phi_2)} \\
 &= a_1^2 + a_2^2 + 2 \cos(\phi_1 - \phi_2) a_1 a_2
 \end{aligned}$$

$$\begin{aligned}
 \cos x &= \frac{e^{+jx} + e^{-jx}}{2} \\
 &= \operatorname{Re} \left\{ e^{jx} \right\}
 \end{aligned}$$

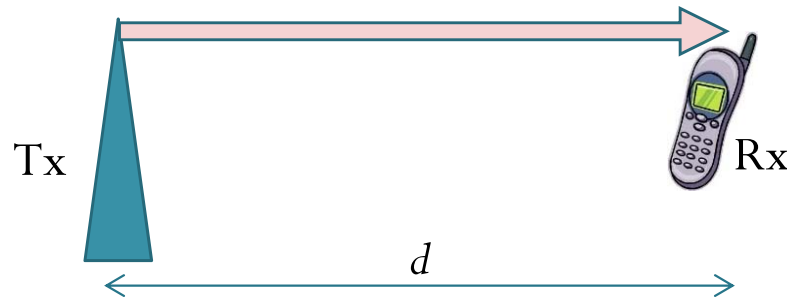
Ray tracing (a revisit)

$$P_r = \left(\frac{\sqrt{G_{Tx} G_{Rx}} \lambda}{4\pi d} \right)^2 P_t = \left(\frac{\alpha}{d} \right)^2 P_t$$

- LOS: $P_r = \frac{(\sqrt{2P_t})^2}{2} = P_t$

$$x(t) = \sqrt{2P_t} \cos(2\pi f_c t)$$

$$y(t) = \frac{\alpha}{d} \sqrt{2P_t} \cos\left(2\pi f_c \left(t - \frac{d}{c}\right)\right)$$



From Friis equation,

$$\alpha = \frac{\sqrt{G_{Tx} G_{Rx}} \lambda}{4\pi}$$

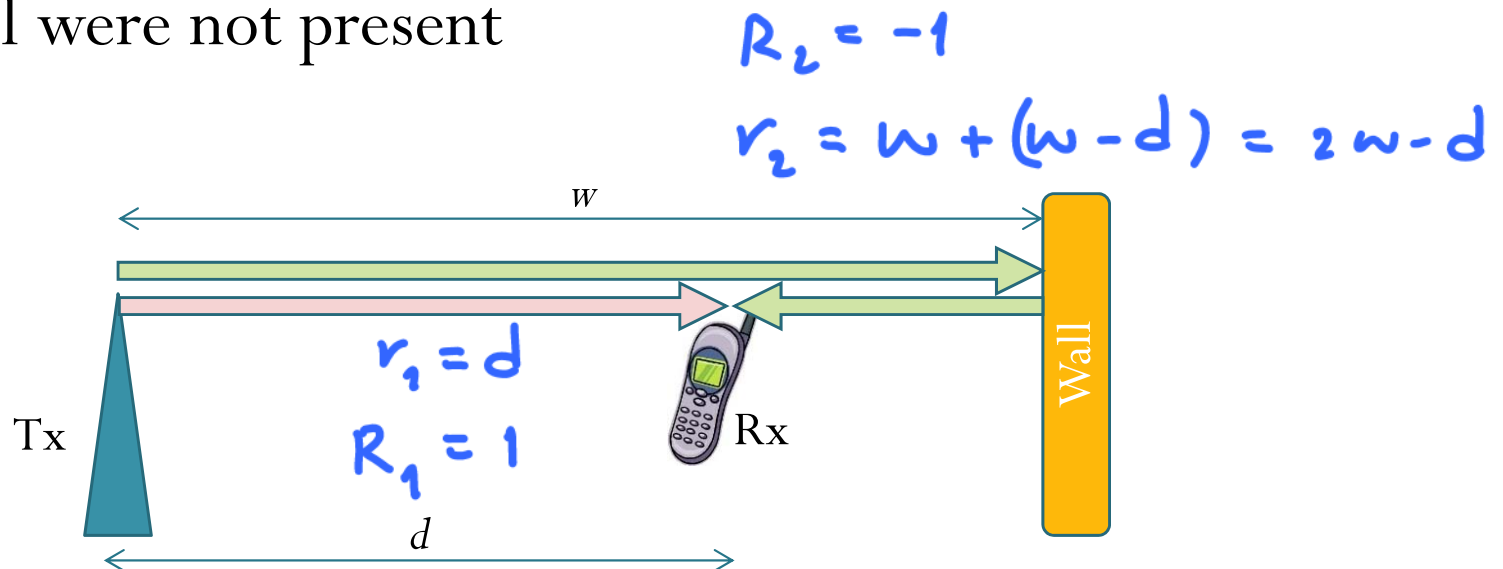
- Multipath Reception

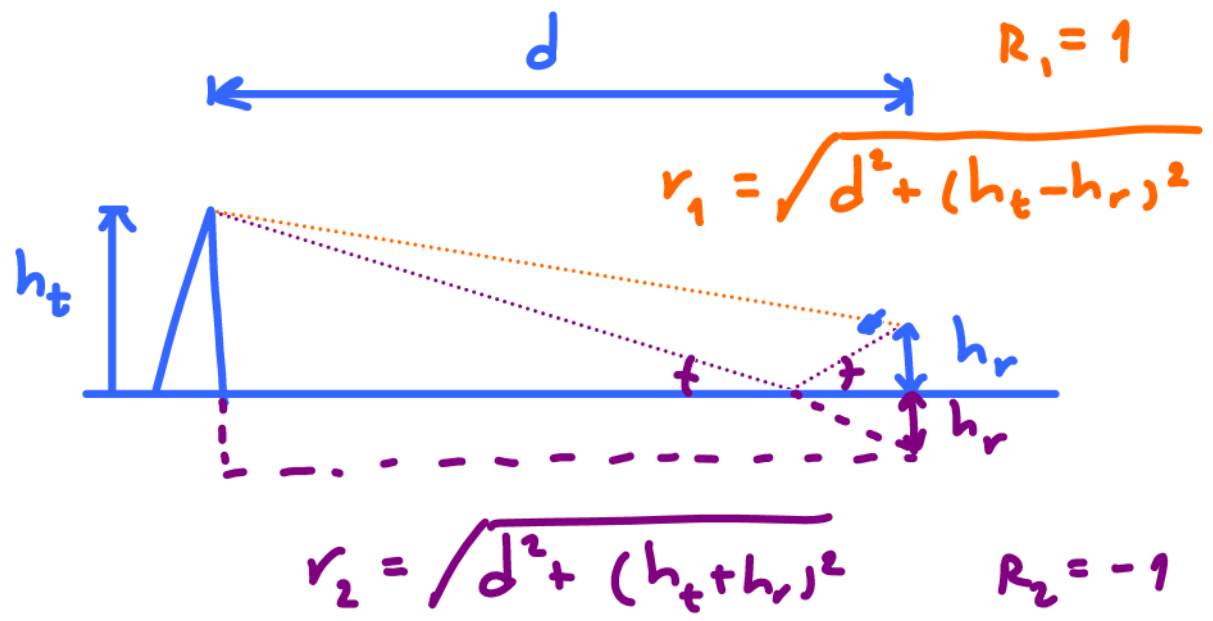
$$y(t) = \sum_{k=1}^n R_k \frac{\alpha}{r_k} \sqrt{2P_t} \cos\left(2\pi f_c \left(t - \frac{r_k}{c}\right)\right)$$

reflection coefficient
 assume $\begin{cases} 1 & \text{for no reflection} \\ -1 & \text{for one reflection} \end{cases}$

Ex. One reflecting wall (1/4)

- There is a fixed antenna transmitting the sinusoid $x(t)$, a fixed receive antenna, and a single perfectly reflecting large fixed wall.
- Assume that the wall is very large, the reflected wave at a given point is the same (except for a sign change) as the free space wave that would exist on the opposite side of the wall if the wall were not present





Ex. One reflecting wall (2/4)

$$x(t) = \sqrt{2P_t} \cos(2\pi f_c t)$$

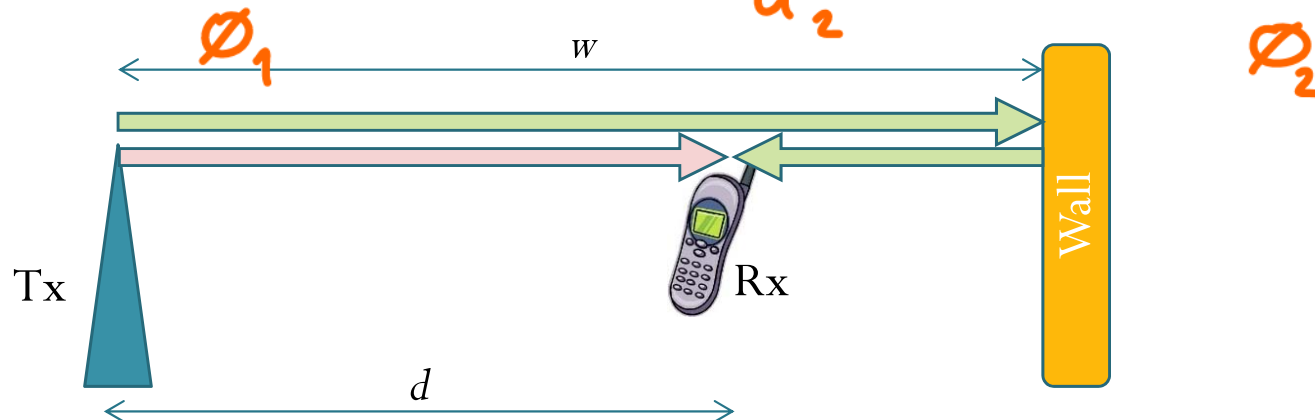
$$y(t) = \frac{\alpha}{d} \sqrt{2P_t} \cos\left(2\pi f_c \left(t - \frac{d}{c}\right)\right) - \frac{\alpha}{2w-d} \sqrt{2P_t} \cos\left(2\pi f_c \left(t - \frac{2w-d}{c}\right)\right)$$

$$= \frac{\alpha}{d} \sqrt{2P_t} \cos\left(2\pi f_c \left(t - \frac{d}{c}\right)\right) - \frac{\alpha}{2w-d} \sqrt{2P_t} \cos\left(2\pi f_c \left(t - \frac{2w-d}{c}\right)\right)$$

$$= \frac{\alpha}{d} \sqrt{2P_t} \cos\left(2\pi f_c \left(t - \frac{d}{c}\right)\right) + \frac{\alpha}{2w-d} \sqrt{2P_t} \cos\left(2\pi f_c \left(t - \frac{2w-d}{c}\right) - \pi\right)$$

a_1

a_2



Ex. One reflecting wall (3/4)

$$y(t) = \frac{\alpha}{d} \sqrt{2P_t} \cos\left(2\pi f_c \left(t - \frac{d}{c}\right)\right) + \frac{\alpha}{2w-d} \sqrt{2P_t} \cos\left(2\pi f_c \left(t - \frac{2w-d}{c}\right) - \pi\right)$$

$$P_y = P_t \left(\left(\frac{\alpha}{d}\right)^2 + \left(\frac{\alpha}{2w-d}\right)^2 + 2 \frac{\alpha^2}{d(2w-d)} \cos(\Delta\phi) \right)$$

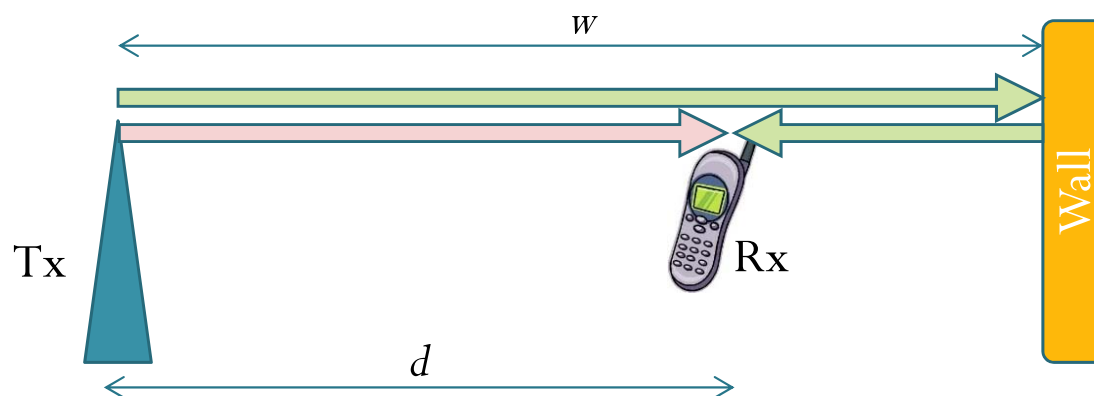
$$\Delta\phi = 2\pi f_c \frac{2w-2d}{c} + \pi = 2\pi \frac{1}{\lambda/2} (w-d) + \pi$$

period = $\frac{\lambda}{2}$

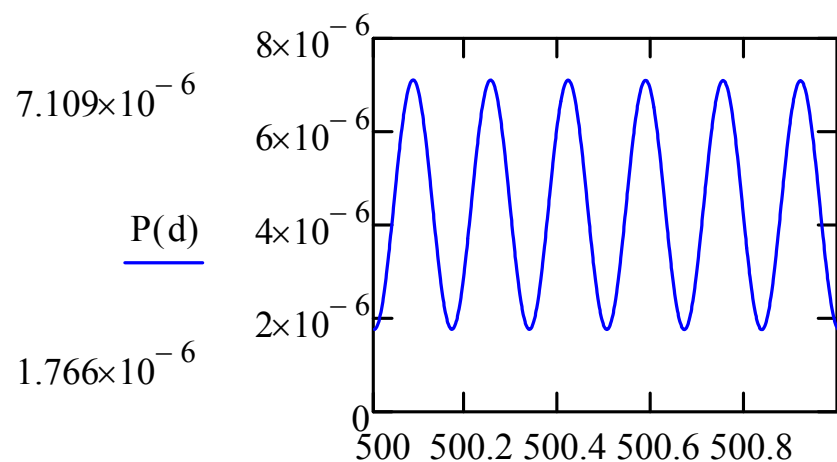
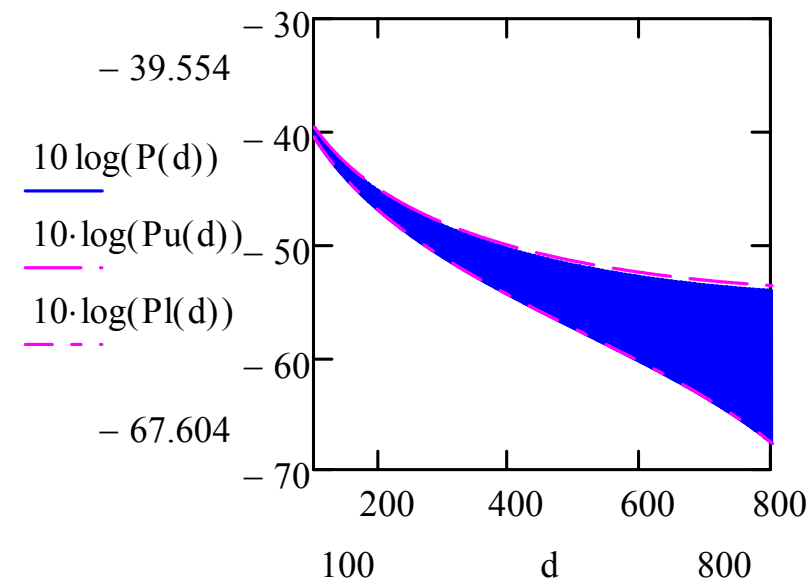
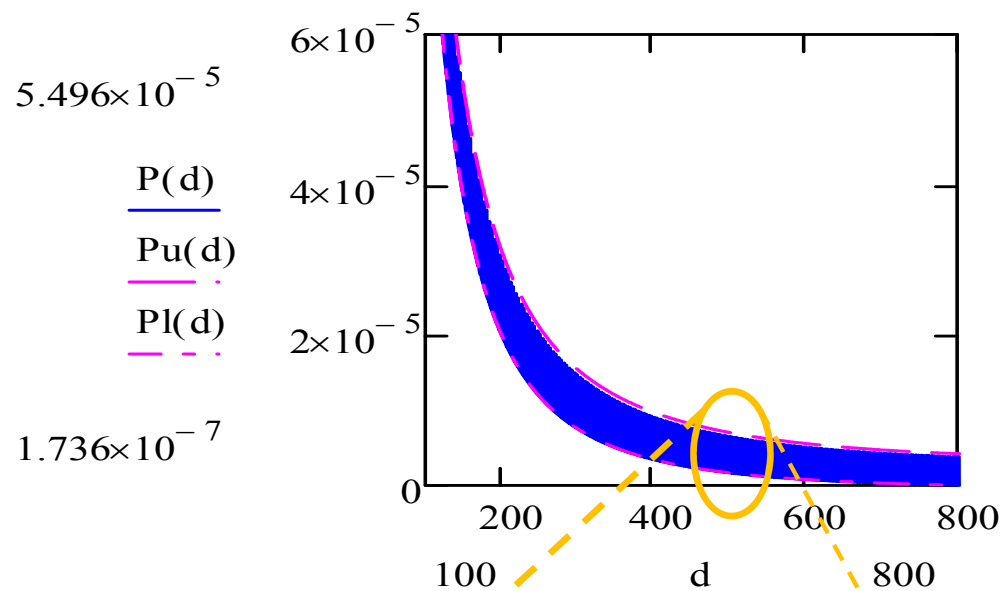
$\cos(2\pi 5x)$



form constructive and destructive interference pattern



Ex. One reflecting wall (4/4)



period = $\frac{1}{6} \approx 0.167 \text{ m}$

$f = 900 \text{ MHz}$

$w = 1 \text{ km}$

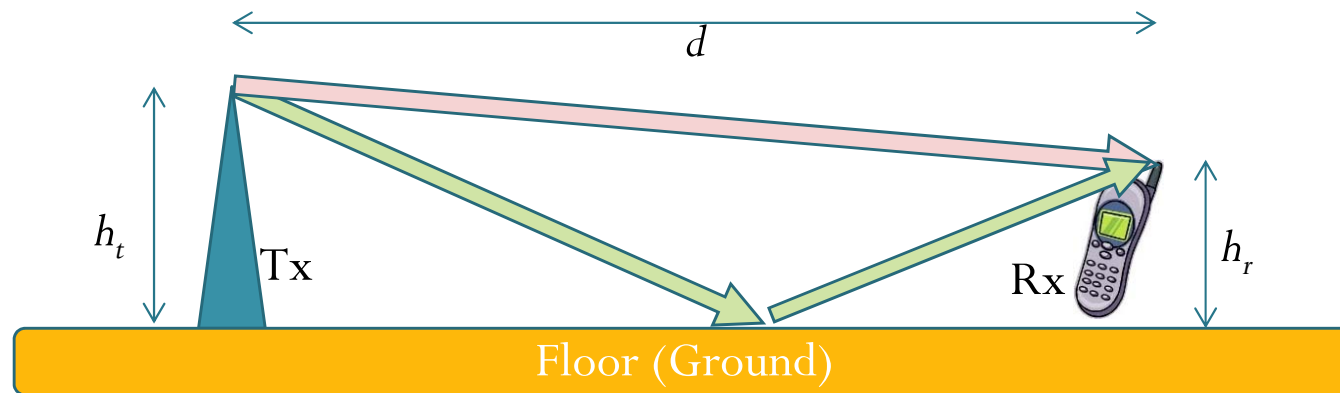
$\lambda = \frac{c}{f} = \frac{3 \times 10^8}{9 \times 10^8} = \frac{1}{3}$

Ex. Two-Ray Model

$$\text{Delay spread} = \frac{r_2}{c} - \frac{r_1}{c}$$

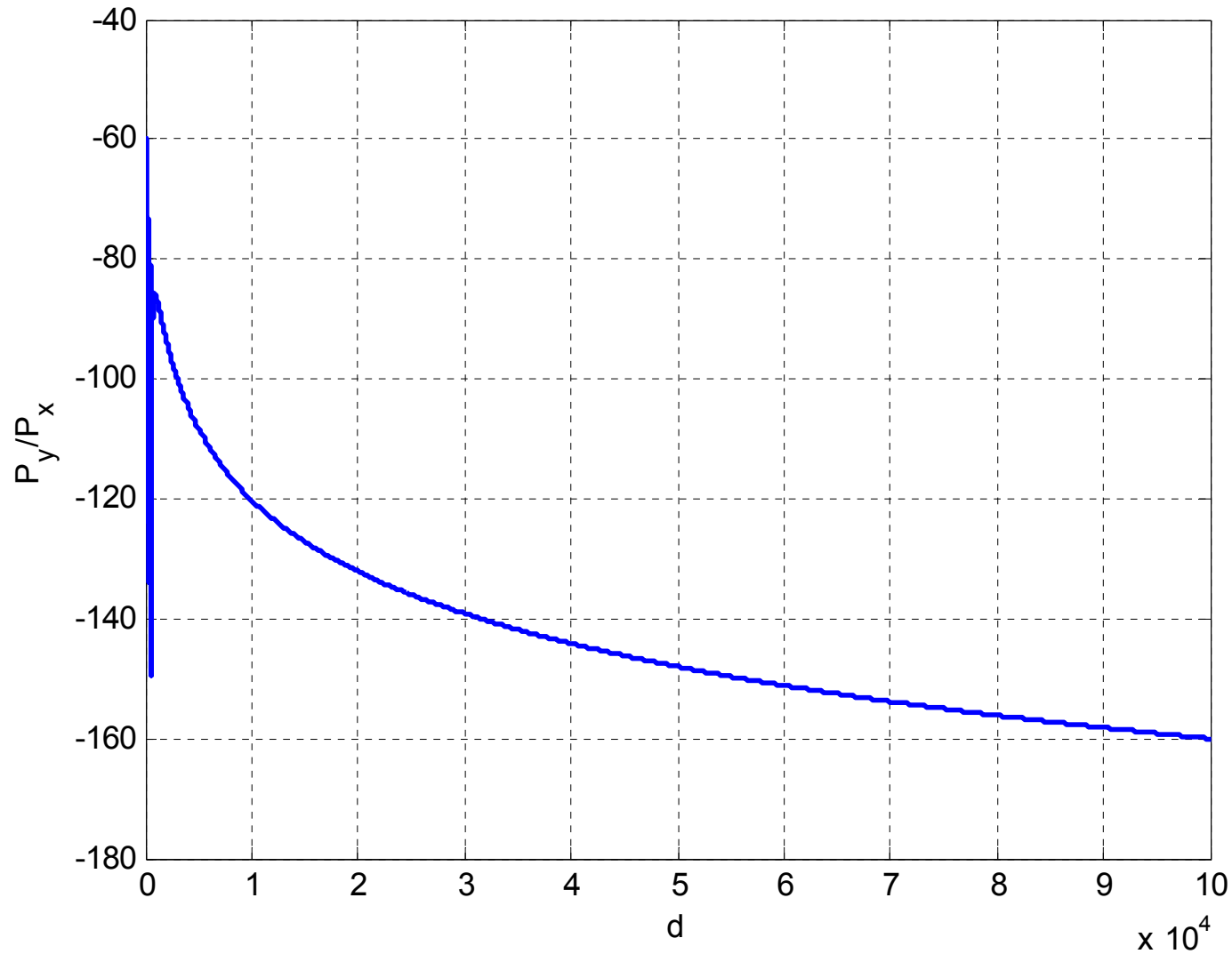
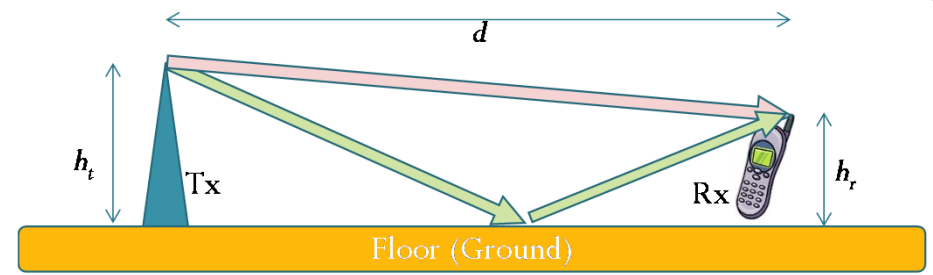
$$y(t) = \frac{\alpha}{r_1} \sqrt{2P_t} \cos\left(2\pi f_c \left(t - \frac{r_1}{c}\right)\right) - \frac{\alpha}{r_2} \sqrt{2P_t} \cos\left(2\pi f_c \left(t - \frac{r_2}{c}\right)\right)$$

$$\frac{P_y}{P_x} = \left| \frac{\alpha}{r_1} e^{-j2\pi f_c \frac{r_1}{c}} - \frac{\alpha}{r_2} e^{-j2\pi f_c \frac{r_2}{c}} \right|^2 = \left| \frac{\alpha}{r_1} - \frac{\alpha}{r_2} e^{-j2\pi f_c \frac{r_2 - r_1}{c}} \right|^2$$



Assume ground reflection coefficient = -1.

Ex. Two-Ray Model

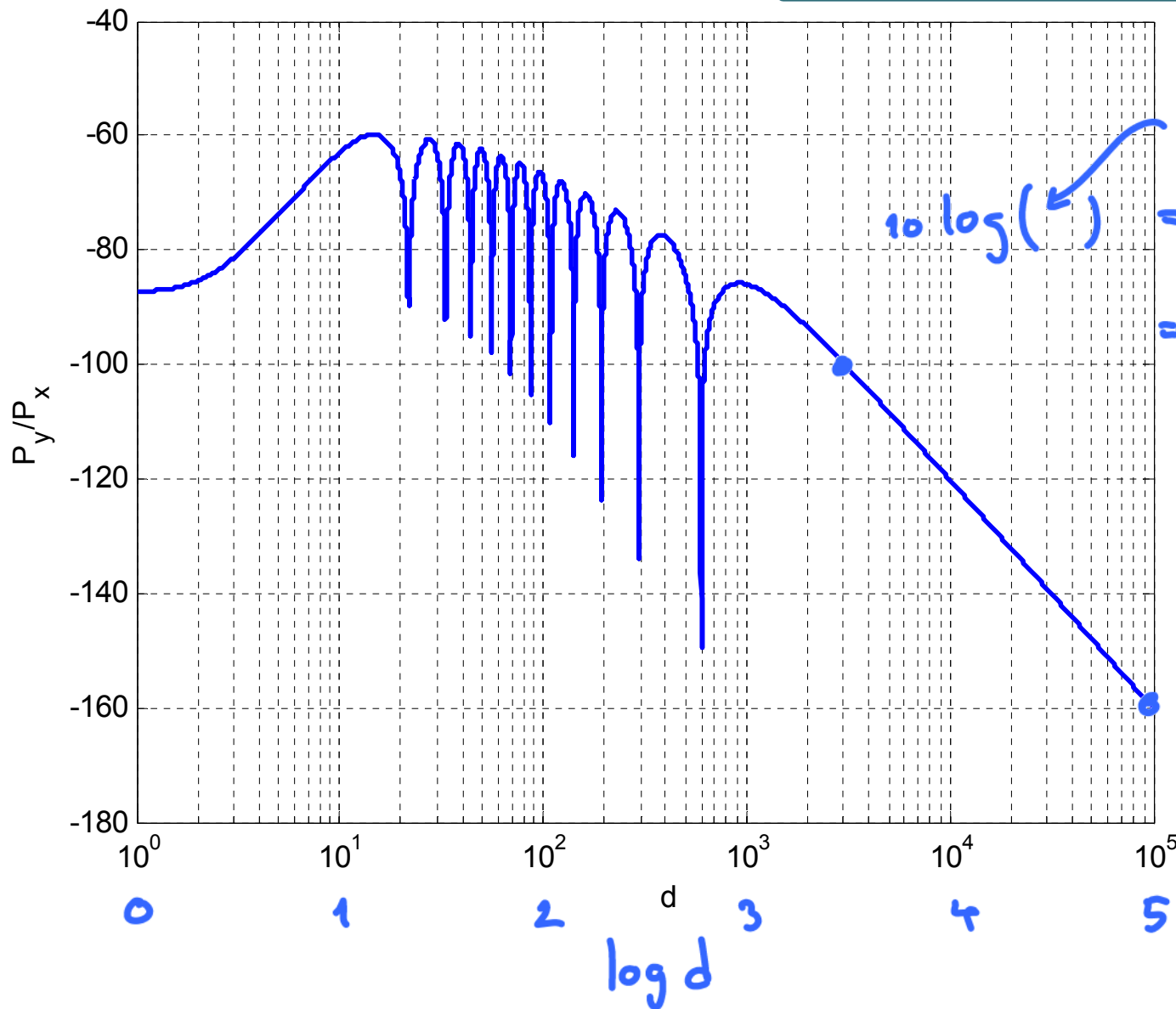
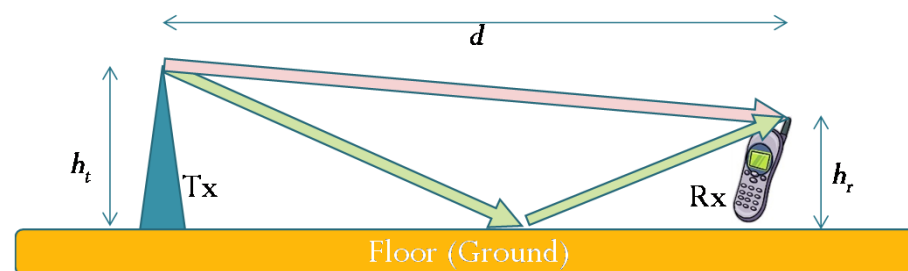


$$f = 900 \text{ MHz}$$

$$h_t = 50 \text{ m}$$

$$h_r = 2 \text{ m}$$

Ex. Two-Ray Model



$$\frac{P_r}{P_t} = K \left(\frac{d_0}{d} \right)^\alpha$$

$$= 10 \log K + \alpha 10 \log \frac{d_0}{d}$$

$$= () - 10\alpha \log d$$

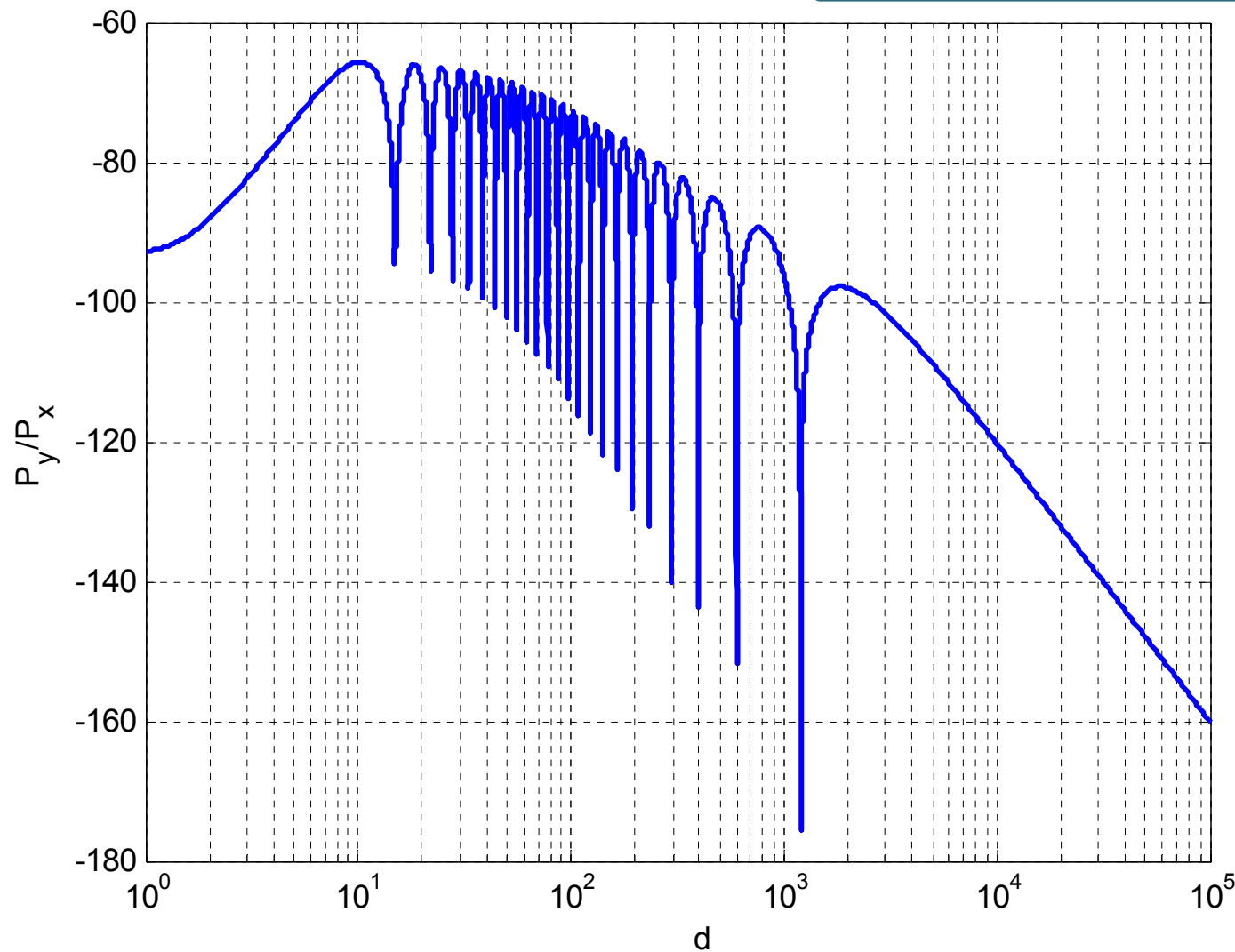
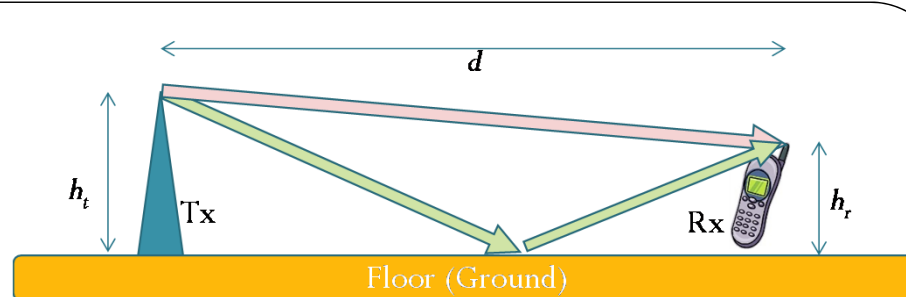
$f = 900 \text{ MHz}$

$h_t = 50 \text{ m}$

$h_r = 2 \text{ m}$

slope ≈ -40

Ex. Two-Ray Model



$$f = 1800 \text{ MHz}$$

$$h_t = 50 \text{ m}$$

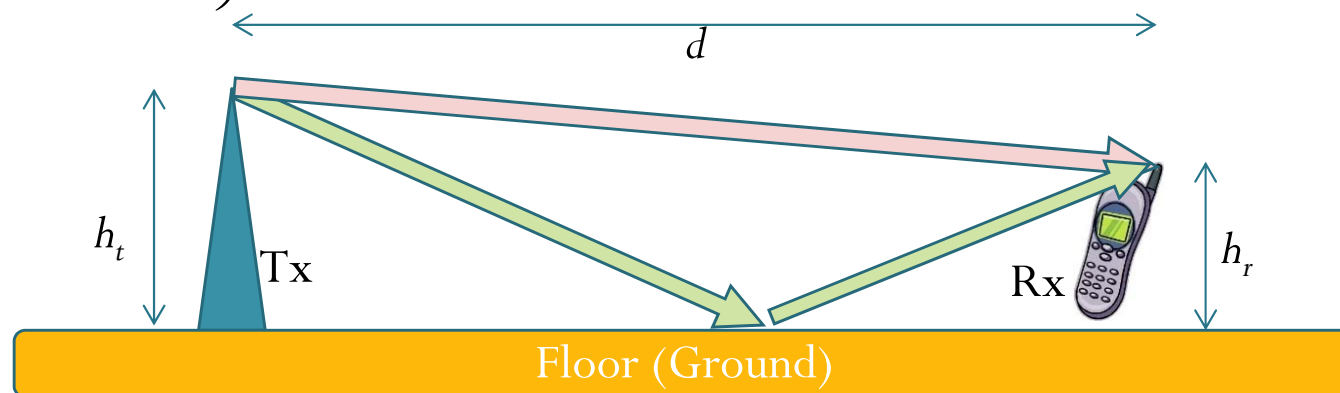
$$h_r = 2 \text{ m}$$

Ex. Two-Ray Model (Approximation)

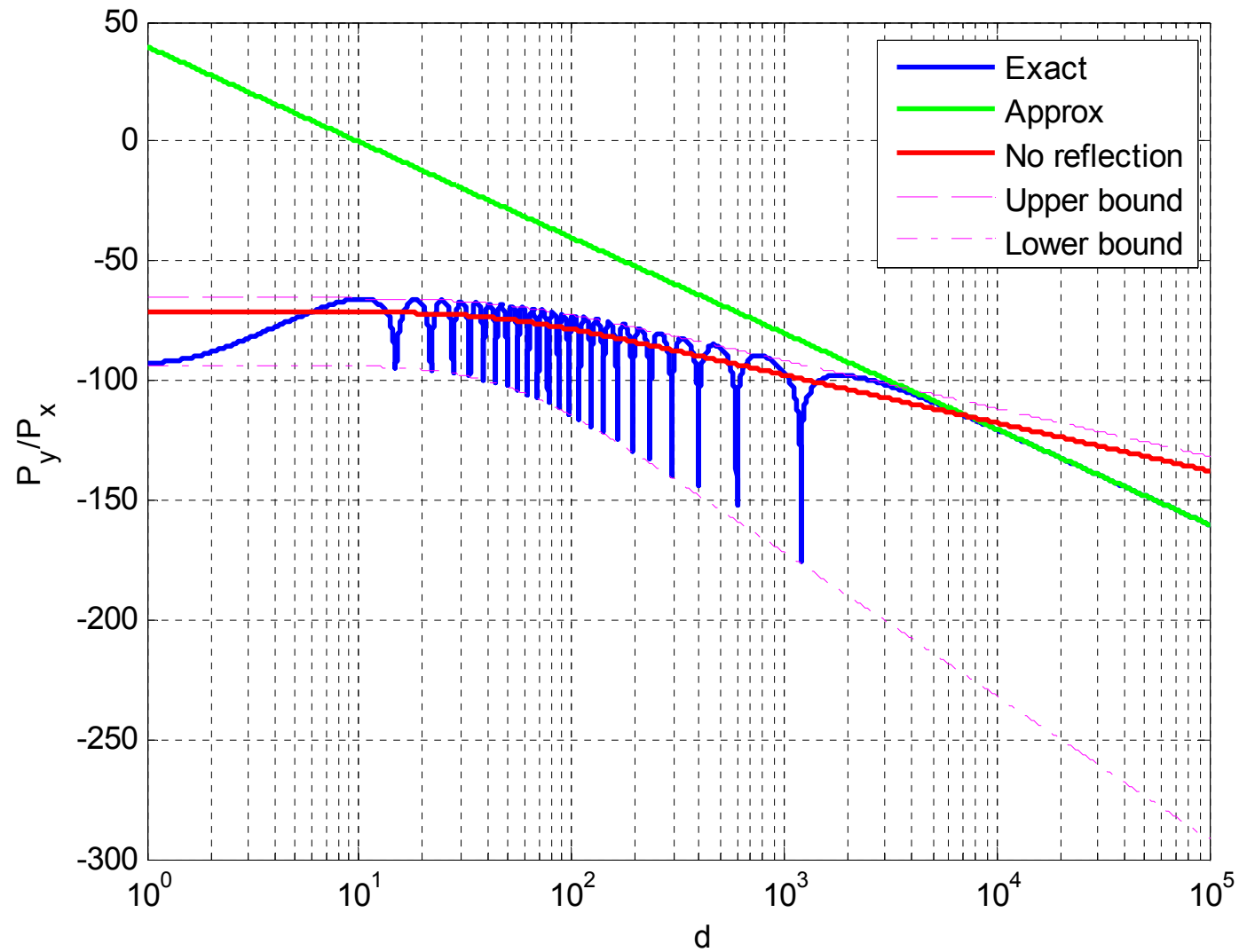
$$\frac{P_y}{P_x} \approx \left| \frac{\alpha}{r_1} - \frac{\alpha}{r_2} e^{-j2\pi \frac{2h_t h_r}{\lambda}} \right|^2 \approx \frac{\alpha}{d} \left| 1 - e^{-j2\pi \frac{2h_t h_r}{\lambda d}} \right|^2 \quad d \gg h_t, h_r$$

$$\approx \left(\frac{\alpha}{d} \right)^2 \left| 1 - \left(1 - j2\pi \frac{2h_t h_r}{\lambda d} \right) \right|^2 = \frac{\alpha^2}{d^2} \left| j2\pi \frac{2h_t h_r}{\lambda d} \right|^2 = \left(\frac{4\pi\alpha h_t h_r}{\lambda d^2} \right)^2$$

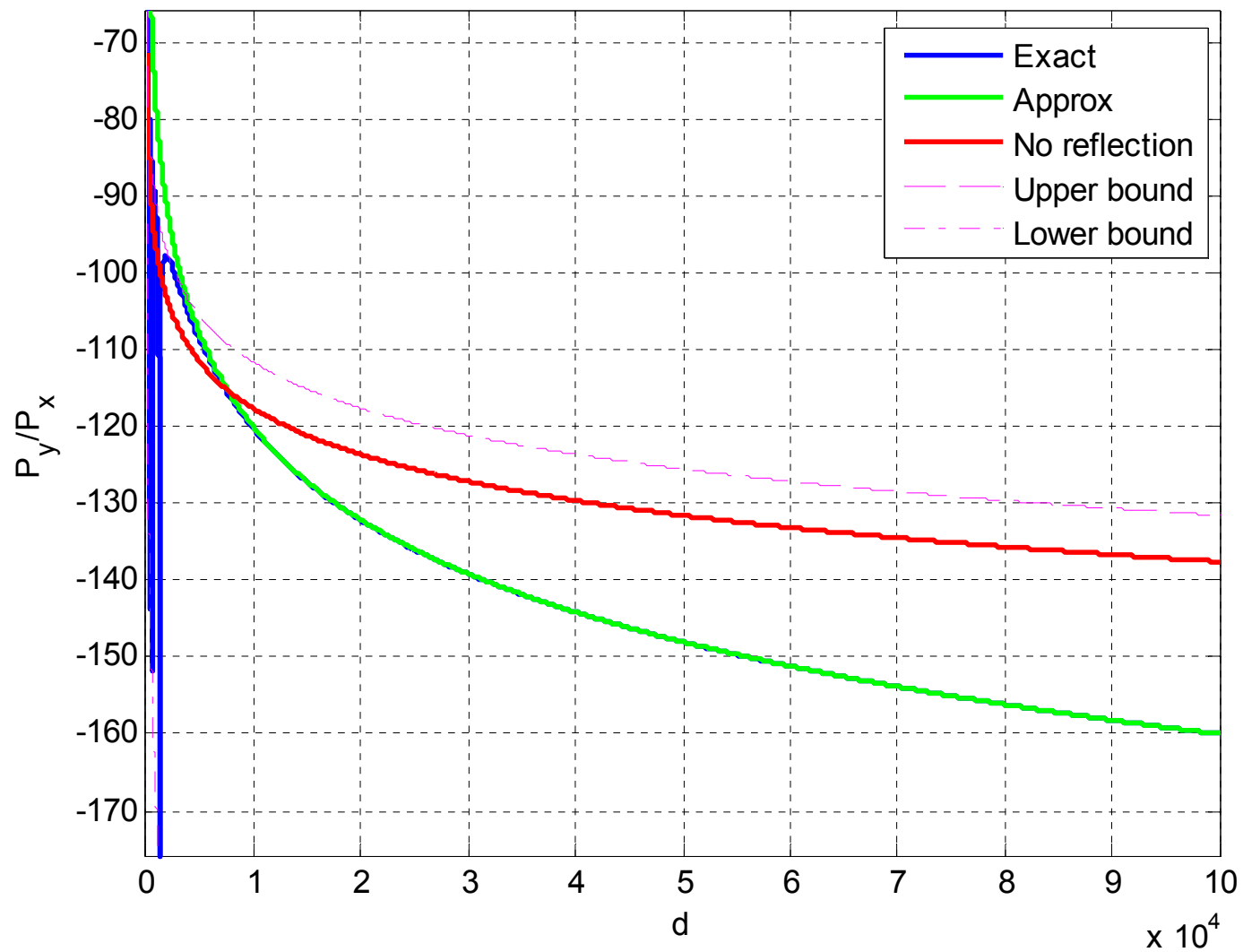
$$= \left(\frac{\sqrt{G_{Tx} G_{Rx}} h_t h_r}{d^2} \right)^2 \propto \frac{1}{d^4}$$



Ex. Two-Ray Model



Ex. Two-Ray Model

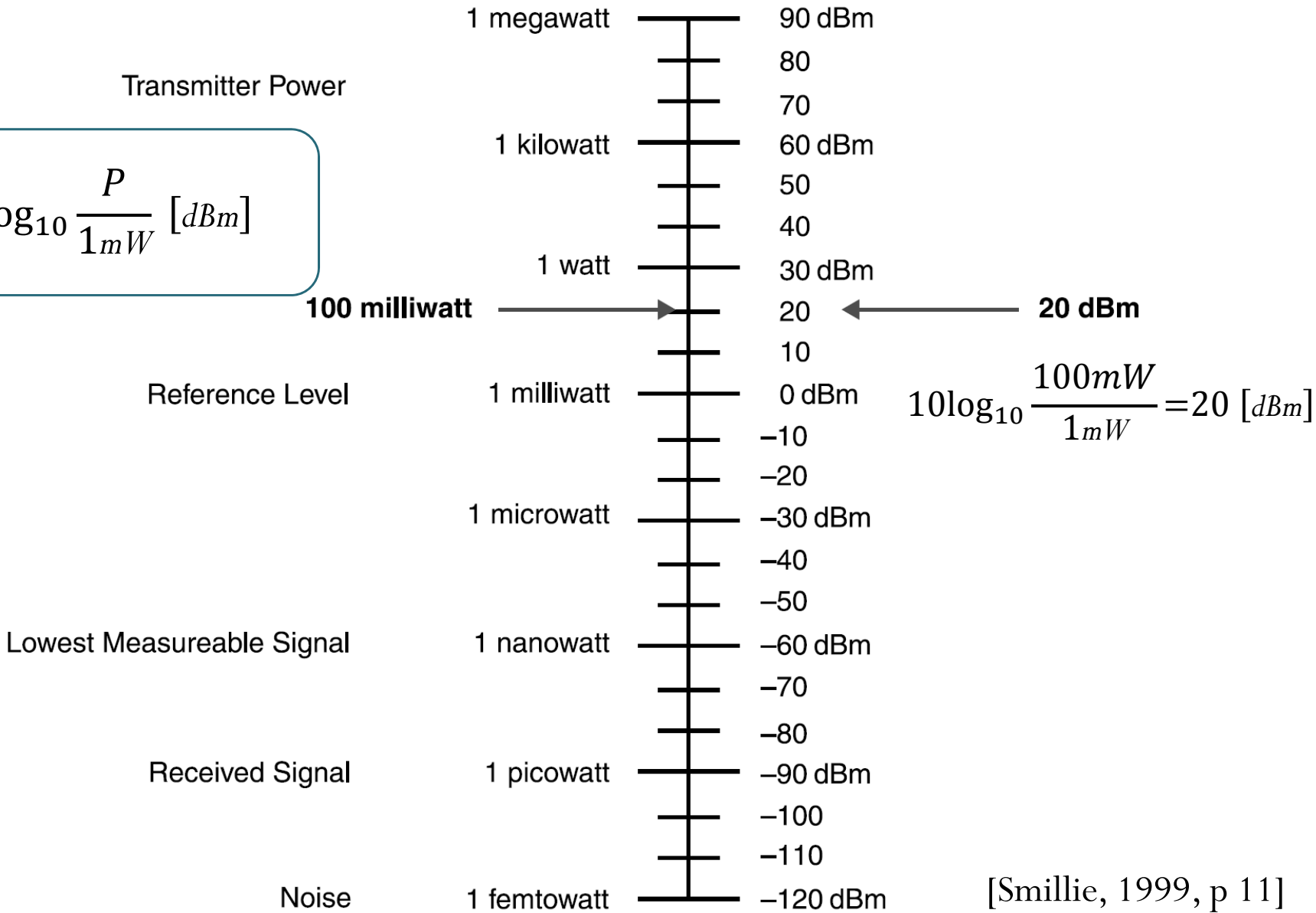


dBm

- The range of RF power that must be measured in cellular phones and wireless data transmission equipment varies from
 - hundreds of watts in base station transmitters to
 - picowatts in receivers.
- For calculations to be made, all powers must be expressed in the same power units, which is usually **milliwatts**.
 - A transmitter power of 100 W is therefore expressed as 100,000mW. A received power level of 1 pW is therefore expressed as 0.000000001mW.
- Making power calculations using decimal arithmetic is therefore complicated.
- To solve this problem, the dBm system is used.

Range of RF Power in Watts and dBm

$$P [W] = 10 \log_{10} \frac{P}{1mW} [dBm]$$



dB and dBm

- The decibel scale expresses factors or ratios logarithmically.
- Unitless dB value
 - Represent power ratio: $10\log_{10} \frac{P_2}{P_1}$
- dB value with a unit

- Represent the signal power itself:

$$P[\text{dBW}] = 10\log_{10} \frac{P}{1 \text{ W}}, \quad P[\text{dBm}] = 10\log_{10} \frac{P}{1 \text{ mW}}$$

- Note that $P[\text{dBm}] = P[\text{dBW}] + 30$

Remark

- Adding dB values corresponds to multiplying the underlying factors, which means multiplying the units if they are present.
- It is therefore appropriate to add unitless dB values to a dB value with a unit (such as dBm)
 - The result is still referred to that unit.
 - Ex: $17 \text{ dBm} + 13 \text{ dB} - 6 \text{ dB} = 24 \text{ dBm}$
 - Correspond to $50 \text{ mW} \times 20 / 4 = 250 \text{ mW}$.

Doppler Shift: 1D Move

- At the transmitter, suppose we have

$$\sqrt{2P_t} \cos(2\pi f_c t + \phi)$$

- At distance r (far enough), we have **Time to travel a distance of r**

$$\frac{\alpha}{r} \sqrt{2P_t} \cos\left(2\pi f_c \left(t - \frac{r}{c}\right) + \phi\right)$$

- If moving, r becomes $r(t)$.
- If moving **away** at a constant velocity v , then $r(t) = r_0 + vt$.

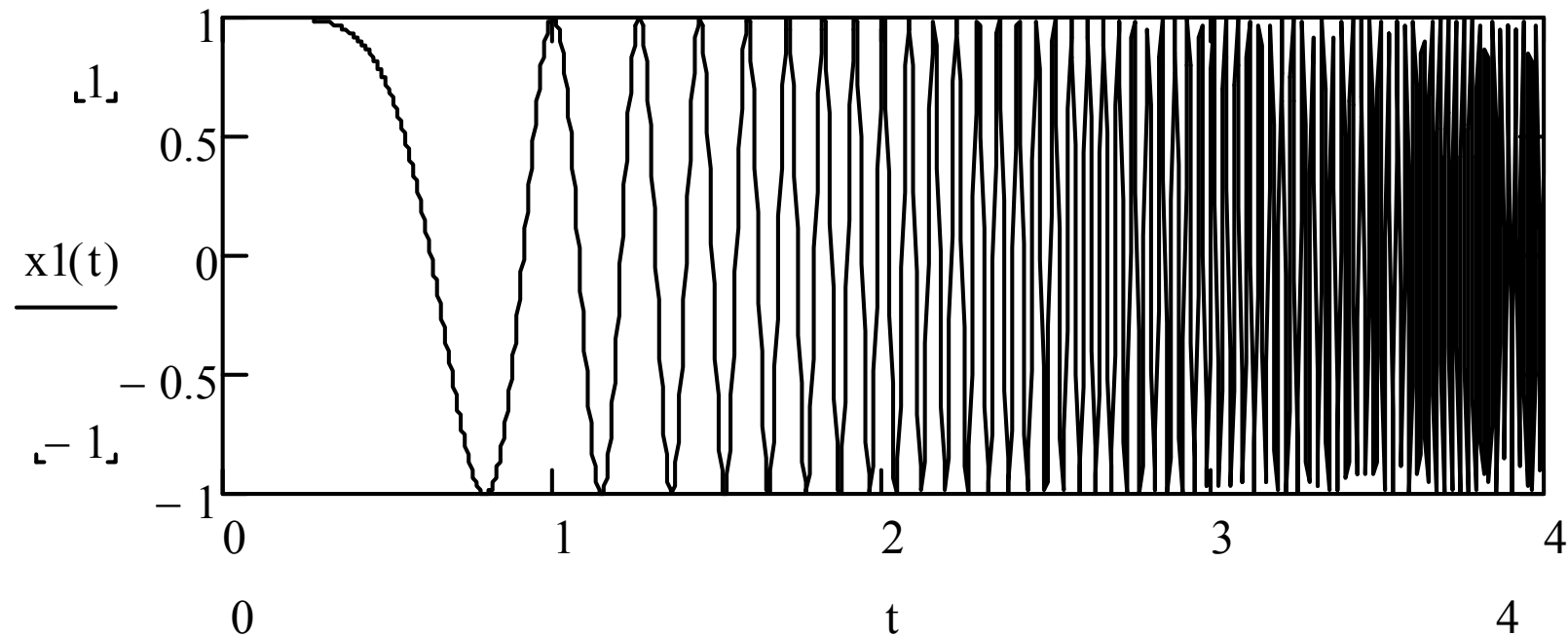
$$\frac{\alpha}{r(t)} \cos\left(2\pi f_c \left(t - \frac{r_0 + vt}{c}\right) + \phi\right) = \frac{\alpha}{r(t)} \cos\left(2\pi \left(f_c - f_c \frac{v}{c}\right) t - 2\pi f_c \frac{r_0}{c} + \phi\right)$$

Frequency shift

$$\Delta f = \frac{v}{\lambda}$$

Instantaneous Frequency

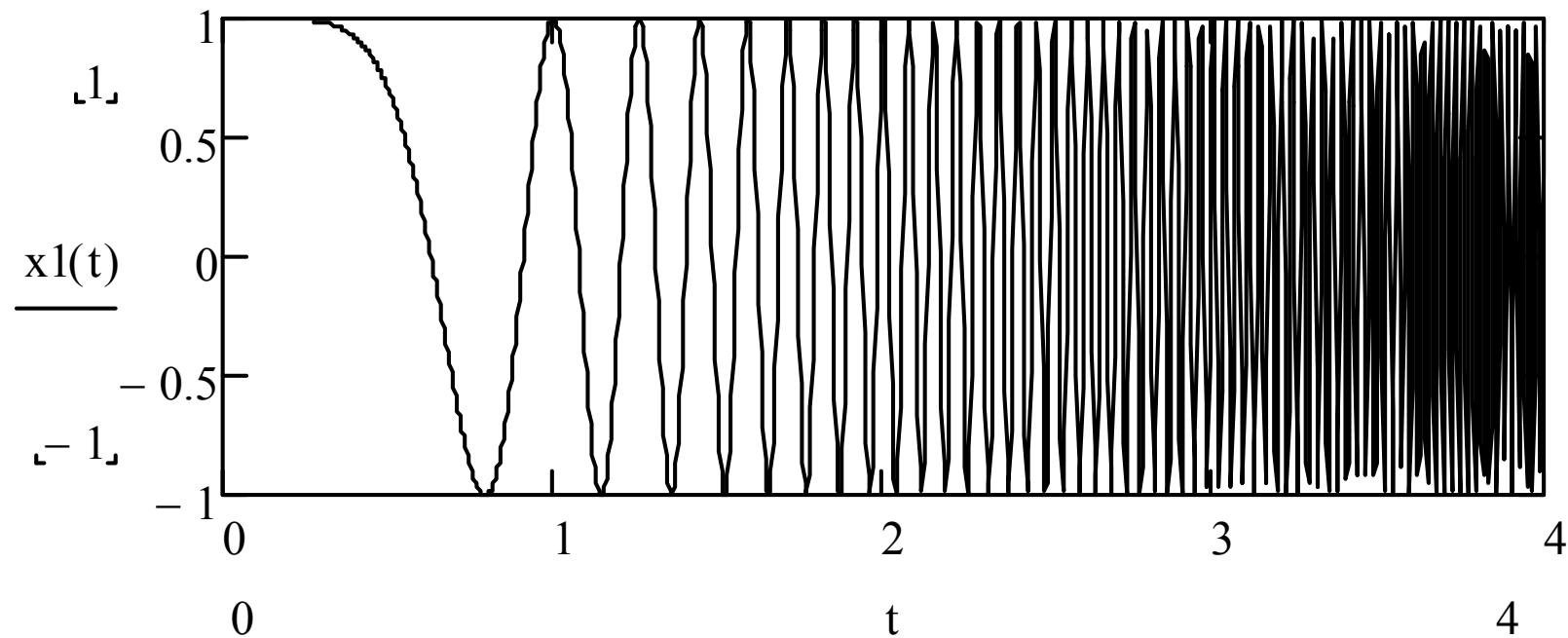
$$x_1(t) = \cos(2\pi t^2 t)$$



At $t = 2$, frequency = ?

Instantaneous Frequency

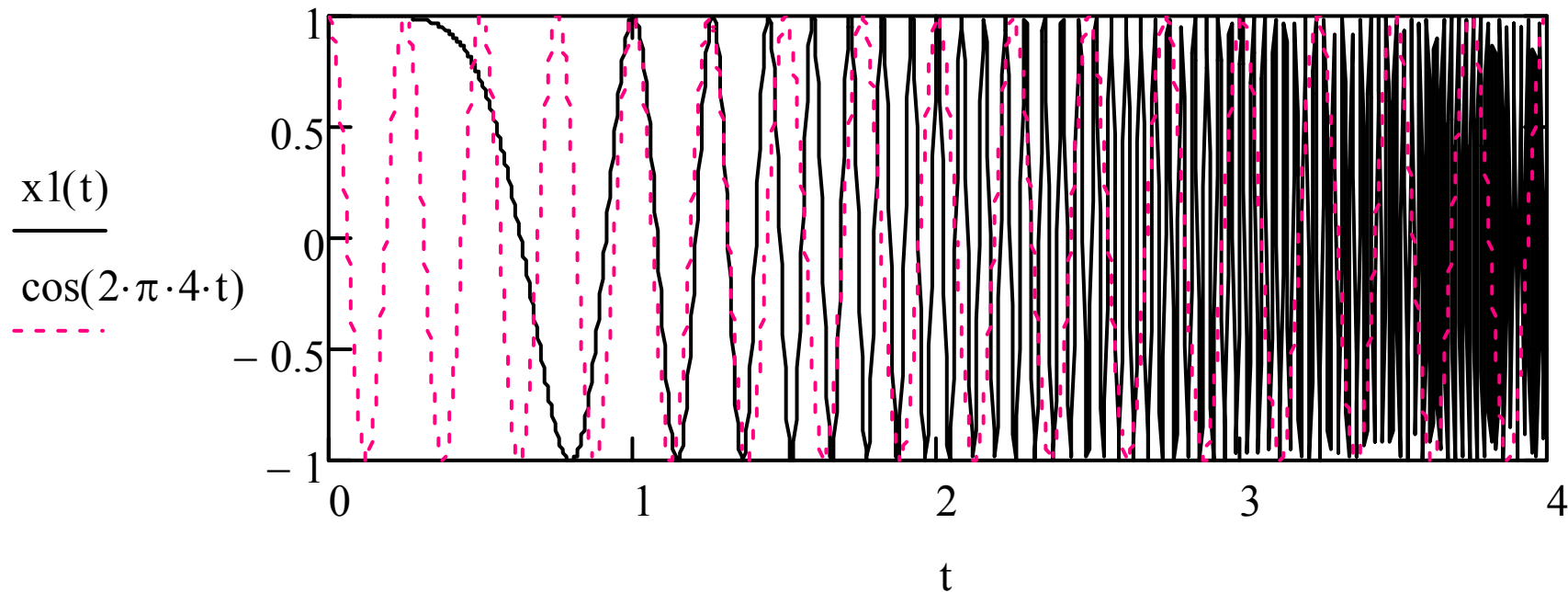
$$x_1(t) = \cos(2\pi t^2 t)$$



$\cos(2\pi ft)$ \longrightarrow At $t = 2$, $f = t^2 = 4$ Hz?

Instantaneous Frequency

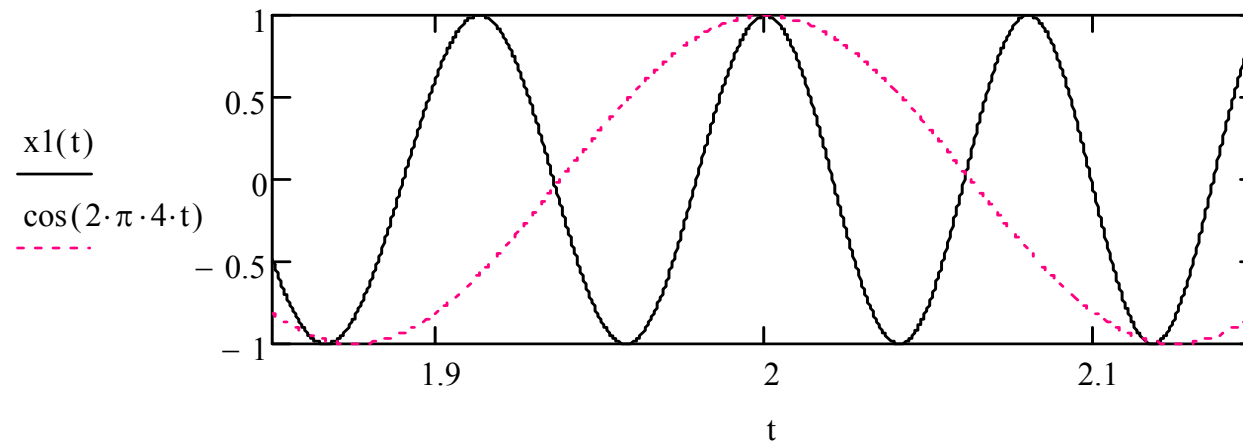
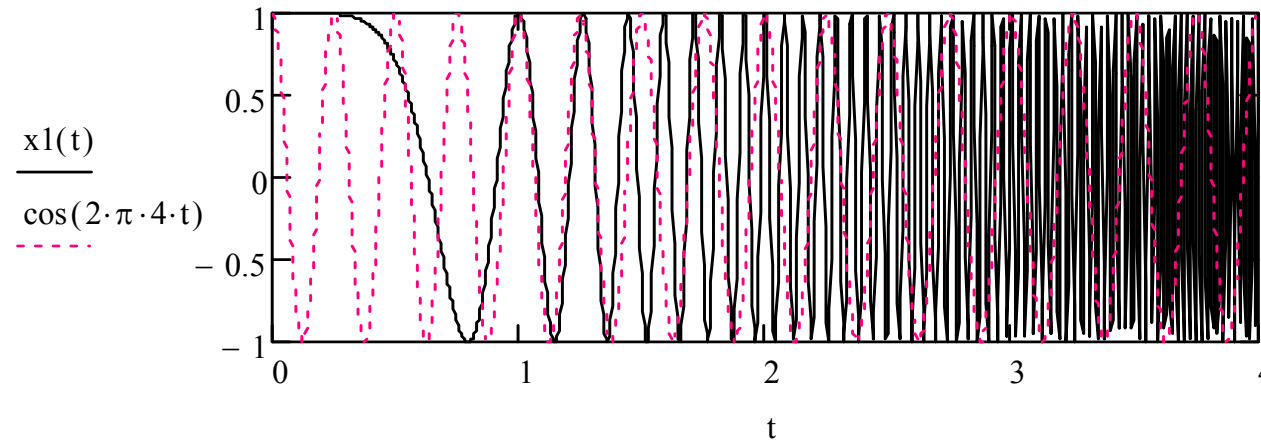
$$x_1(t) = \cos(2\pi t^2 t)$$



$\cos(2\pi ft)$ \longrightarrow At $t = 2$, $f = t^2 = 4$ Hz?

Instantaneous Frequency

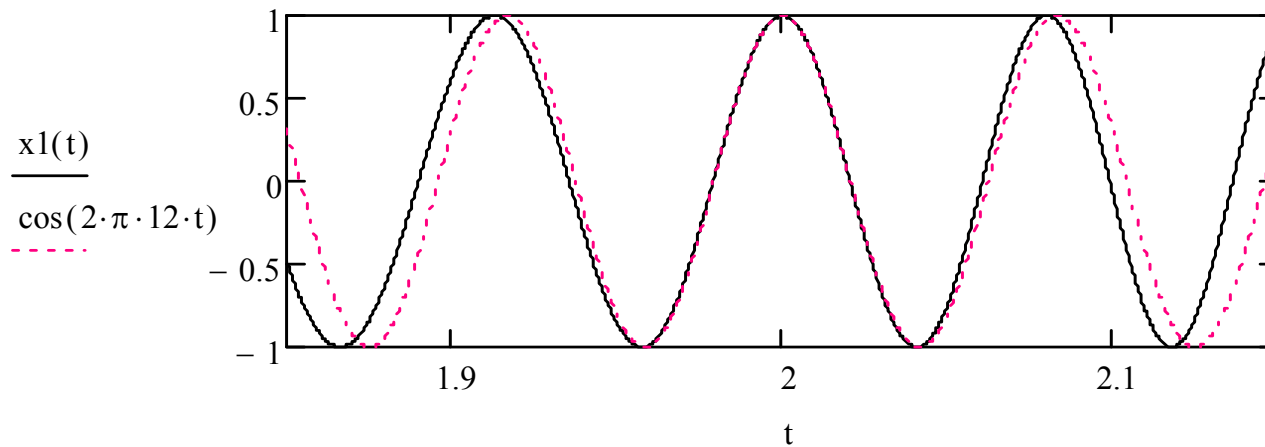
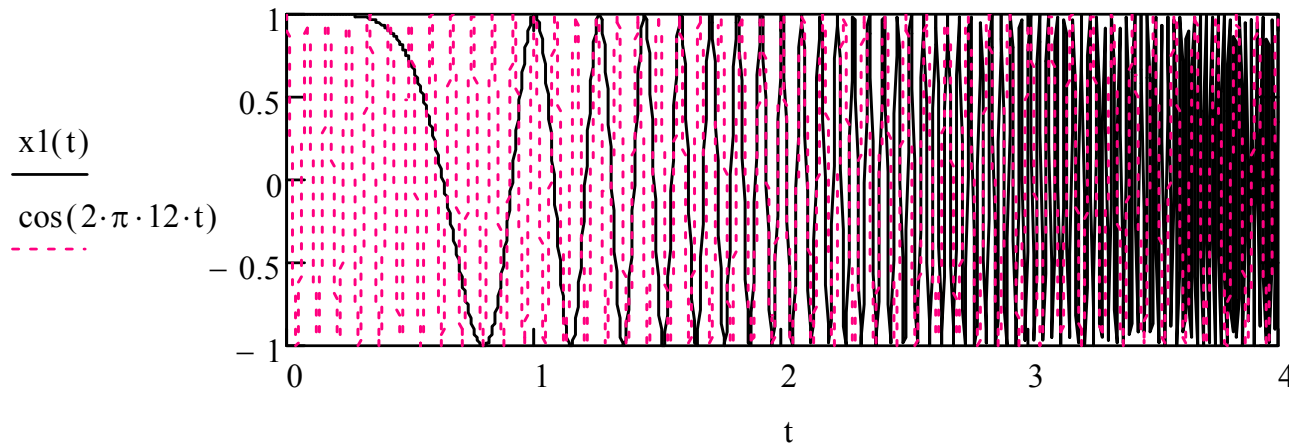
$$x_1(t) = \cos(2\pi t^2 t)$$



4 Hz is too low!!!

Instantaneous Frequency

$$x_1(t) = \cos(2\pi t^2 t)$$



12 Hz?

Review: Instantaneous Frequency

For a generalized sinusoid signal

$$A \cos(\theta(t)),$$

the **instantaneous frequency** at time t is given by

$$f(t) = \frac{1}{2\pi} \frac{d}{dt} \theta(t).$$

When $\theta(t) = 2\pi f_c \left(t - \frac{r(t)}{c} \right) + \phi,$

$$f(t) = \frac{1}{2\pi} \frac{d}{dt} \theta(t) = f_c - \frac{f_c}{c} \frac{d}{dt} r(t) = f_c - \frac{1}{\lambda} \frac{d}{dt} r(t)$$

Frequency shift

Big Picture

Transmission impairments in cellular systems

Physics of radio propagation

Attenuation (Path Loss)

Shadowing

Doppler shift

Inter-symbol interference (ISI)

Flat fading

Frequency-selective fading

Extraneous signals

Co-channel interference

Adjacent channel interference

Impulse noise

White noise

Transmitting and receiving equipment

White noise

Nonlinear distortion

Frequency and phase offset

Timing errors